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ABSTRACT

The nonlinear dielectric response (NDR) of $Bi_{3.25}La_{0.75}Ti_3O_{12}-Al_2O_3$ systems has been studied. To reveal the effects of the dielectric layer (DL) systematically, a method based on the equivalent circuit analysis and parametric study is proposed. The consistency between the calculated results and the experimental results manifests the validity of the method. In addition to the Rayleigh coefficients varying monotonically with the capacitance of DL, it is revealed that the slope between the real and the imaginary parts of the complex capacitance under different ac voltages is related to the relative loss values of DL and the ferroelectric layer. The frequency dependence of the NDR parameters is also discussed. The method established in this work is useful for research on the devices based on the ferroelectric-dielectric system and the ultra-thin ferroelectric films.

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Ferroelectric materials have attracted tremendous attention due to their potential applications in nonvolatile memories, artificial synapses, and energy storage.^{1–4} The ferroelectric performances originate from the development of domains and the displacement of domain walls under the electric field. The nonlinear dielectric response (NDR) has been used to characterize domain-related processes since Taylor and co-workers proved the contribution of domain walls to the permittivity of Pb(Zr_xTi_{1-x})O₃ (PZT).⁵⁶ According to NDR, the relationship between the real part (ε') of the complex permittivity and the ac electric field is given by the Rayleigh law,⁶

$$\varepsilon'(E_0) = \varepsilon'_{in} + \alpha E_0, \tag{1}$$

where E_0 is the amplitude of the ac field. The initial permittivity ε'_{in} represents the reversible component in the permittivity, and the Rayleigh coefficient α represents the irreversible contribution from the movement of domain walls. Moreover, the imaginary part (ε'') of the complex permittivity also changes with the ac field and follows:⁷

$$\varepsilon''(E_0) = 4\alpha E_0/3\pi.$$
 (2)

From (1) and (2), the relationship between ε' and ε'' can be rewritten as

$$\varepsilon' = \varepsilon'_{in} + 3\pi\varepsilon''/4. \tag{3}$$

It means that ε' changes linearly with ε'' , and the slope *k* should be $3\pi/4$. However, the reported values of α are quite different, and deviations of *k* from $3\pi/4$ have been observed in many ferroelectrics.^{7–15} Some mechanisms, such as the electrode effect and the orientation of films, have been proposed to explain these phenomena.⁵

For devices based on the ferroelectric films, the non-ferroelectric interfacial layer between the ferroelectric layer (FL) and the electrodes, which is the so-called dead layer, has a significant impact on the performances of the devices, such as the reduction of the permittivity and the polarization.^{16,17} Alternatively, novel phenomena, such as the negative capacitance effect and the excellent energy storage performances, are observed in the ferroelectric–dielectric (FE–DE) composite system.^{2,18,19} Therefore, studies on the FE–DE composite system.^{2,18,19} Therefore, studies on the FE–DE composite system will reveal the mechanism and extend the applications of the ferroelectric devices. However, few research works on NDR in the FE–DE system were reported. In addition, the dielectric layer (DL) is always considered as an ideal capacitor, and the simple model of the series capacitor is often used in most published literature.^{16,17} The influence of the dielectric loss cannot be considered by this model. In this work, the artificial FE–DE structures formed by the Bi_{3,25}La_{0,75}Ti₃O₁₂ film and



FIG. 1. The ac voltage dependence of the capacitance for (a) BLT, (b) BLT–AO₁₀₀, and (c) BLT–AO₅₀₀. (d) The frequency dependence of $\varepsilon'_{h, Total}$ and α_{Total} .

 Al_2O_3 films are fabricated, and the NDR is studied experimentally. A method based on the equivalent circuit analysis and parametric study is established to reveal the effects of DL on the NDR parameters.

The Bi3.25La0.75Ti3O12 films were prepared on Pt/Ti/SiO2/Si substrates by the sol-gel method and then covered by the amorphous Al₂O₃ films fabricated by atomic layer deposition (ALD) technology with 100 and 500 ALD cycles. The samples were denoted as BLT, BLT-AO₁₀₀, and BLT-AO₅₀₀. Furthermore, Al₂O₃ films with 100 and 500 ALD cycles were also deposited on Pt/Ti/SiO2/Si substrates and denoted as AO100 and AO₅₀₀, respectively. The thicknesses of BLT, AO₁₀₀, and AO₅₀₀ were 320, 9, and 45 nm, respectively, obtained by the spectroscopic ellipsometry. Au top electrodes with the diameter of 200 μ m were sputtered for the electrical measurements. The dielectric responses of the samples were measured by a dielectric analyzer (Alpha-A, Novocontrol Technologies) in the ac voltage range 0.5-2 V and the frequency range 1-100 kHz. Both the parameters of NDR, including the Rayleigh coefficient α , the slope k, and the initial permittivity ε'_{in} can be deduced from the dependence of the complex capacitance on the oscillating voltage (C-OSC), and the geometry size of the samples were taken into account.

As shown in Figs. 1(a)-1(c), the real part (C') of the complex capacitance increases approximately linearly with the amplitude of the ac voltage and decreases with the frequency for the samples. Figure 1(d) exhibits the frequency dependence of the initial permittivity and the Rayleigh coefficient obtained by fitting the curves in Figs. 1(a)-1(c) using Eq. (1). The derived initial permittivity and Rayleigh coefficient

are the total equivalent values of the stack. They are denoted as $\varepsilon'_{in,Total}$ and α_{Total} . Both $\varepsilon'_{in,Total}$ and α_{Total} decrease linearly with the logarithm of the frequency. Similar phenomena have also been observed in other ferroelectrics such as Pb(Yb_{1/2}Nb_{1/2})O₃-PbTiO₃.⁶ Nattermann *et al.* have theoretically presented the dependence of the susceptibility on the logarithm of the frequency in disorder system and explained the pinning of the magnetic domain wall.²⁰ Thereafter, this mechanism is used to explain the pinning of the ferroelectric domain wall.^{21–23} The frequency dependence of $\varepsilon'_{in,Total}$ and α_{Total} can be depicted as

$$\begin{aligned} \varepsilon'_{in,Total} &= e_0 - e \ln(\omega), \\ \alpha_{Total} &= a_0 - a \ln(\omega). \end{aligned} \tag{4}$$

The fitting results of the curves in Fig. 1(d) by Eq. (4) are exhibited in Table I.

It is manifested from Fig. 1(d) that both $\varepsilon'_{in,Total}$ and α_{Total} decrease with the Al₂O₃ thickness increasing. The Rayleigh coefficient is generally considered to be related to the distribution of the potential barriers that hinder the movement of domain walls in ferroelectrics.^{24–26} It is decided by the ferroelectrics. However, when the Rayleigh coefficient is derived from the experimentally measured capacitance, the obtained value is influenced by the existence of the non-ferroelectric interfacial layer, which results in the deviation of the measured capacitance from that of the pure ferroelectric capacitor, just like those demonstrated in this work.

To further investigate the influence of DL on NDR, the relationship between the real part and the imaginary part (C'') of the complex capacitance is considered. According to Eq. (3), the relationship between C' and C'' follows:

$$C' = C'_{in} + kC'',$$
 (5)

and *k* should be $3\pi/4$. Figures 2(a)–2(c) display the curves in the complex capacitance plane. Both *C'* and *C''* decrease when the thickness of the Al₂O₃ film increases. This feature accords with the Kramers–Kronig relations.²⁷ The linear relationship between *C'* and *C''* for the samples can be observed. The fitted initial capacitance C'_{in} , T_{otal} and k_{Total} for the total stack by Eq. (5) is shown in Fig. 2(d). The variations of $C'_{in,Total}$ with frequency and DL thickness are consistent with those of $\varepsilon'_{in,total}$ in Fig. 1(d). Moreover, k_{Total} decreases with the thickness of DL increasing and is larger than the theoretical value of $3\pi/4$. The experimental results indicate that the existence of DL also influences k_{Total} . Because k_{Total} relates to the imaginary part of the complex capacitance and cannot be deduced by the series capacitor model, a method based on the equivalent circuit model and the parametric study is proposed to investigate the effect of DL.

The total complex capacitance (C^*_{Total}) of a circuit composed of two components connected in series is

Sample ID	e ₀ (nF/m)	e (nF s/m rad)	(10^{-18} F/V)	$(10^{-18} \text{ F s/V rad})$
BLT	2.028 ± 0.001	0.0184 ± 0.0001	15.104 ± 0.004	0.498 ± 0.004
BLT-AO ₁₀₀ (Exp.)	1.577 ± 0.001	0.0188 ± 0.0001	6.471 ± 0.003	0.212 ± 0.004
BLT-AO ₁₀₀ (Cal.)	1.3840 ± 0.0003	0.02822 ± 0.00005	6.165 ± 0.003	0.267 ± 0.003
BLT-AO ₅₀₀ (Exp.)	0.6678 ± 0.0004	0.00593 ± 0.00005	0.491 ± 0.003	0.014 ± 0.002
BLT-AO ₅₀₀ (Cal.)	0.6462 ± 0.0002	0.01531 ± 0.00003	1.268 ± 0.001	0.058 ± 0.001

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TABLE I. Fitting results of the samples.

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FIG. 2. *C'* vs *C''* for (a) BLT, (b) BLT–AO₁₀₀, and (c) BLT–AO₅₀₀. (d) The frequency dependence of $C'_{in, Total}$ and k_{Total} . The dash line in (d) is the position of $3\pi/4$. The absence of the data for BLT–AO₅₀₀ in the lower frequency region is because the excessive noise hinders the fitting of the experimental data.

$$C^*_{Total} = 1/(j\omega Z^*_{Total}) = 1/[j\omega (Z^*_{FL} + Z^*_{DL})],$$
(6)

where Z_{Total}^* is the total complex impedance of the circuit, Z_{FL}^* and Z_{DL}^* are the complex impedance of FL and DL, respectively, and ω is the circular frequency. The complex capacitance of FL follows Eq. (3) and can be expressed as



FIG. 3. Relationship between C'_{Total} and G''_{total} with different (a) C'_{DL} and (b) tan δ_{DL} . Mapping diagrams of (c) $C'_{In,Total}$ and (d) k_{Total} .

$$C'_{FL} = C'_{in,FL} + k_{FL}C''_{FL},$$
(7)

where C'_{FL} and C''_{FL} are the real and imaginary parts of the complex capacitance of FL, respectively, $C'_{in,FL}$ and k_{FL} are the initial capacitance and the slope of FL, respectively. Based on Eqs. (6) and (7), C^*_{Total} is rewritten as

$$C_{Total}^{*} = C_{Total}' - jC_{Total}'' = \frac{C_{in,FL}''C_{DL}' + 2C_{in,FL}'C_{DL}''k_{FL}C_{FL}'' + C_{FL}''\left\{C_{DL}'^{2} + C_{DL}''\left[C_{DL}'' + C_{FL}''\left(1 + k_{FL}^{2}\right)\right]\right\}}{(C_{DL}'' + C_{FL}'')^{2} + (C_{in,FL}' + C_{DL} + k_{FL}C_{FL}'')^{2}} - j\frac{C_{in,FL}'^{2}C_{DL}' + C_{in,FL}'\left(C_{DL}^{2} + C_{DL}'' + 2k_{FL}C_{DL}''\right) + C_{FL}''\left[k_{FL}\left(C_{DL}'^{2} + C_{DL}''\right) + C_{DL}'C_{FL}''\left(1 + k_{FL}^{2}\right)\right]}{(C_{DL}'' + C_{FL}'')^{2} + (C_{in,FL}' + C_{DL}' + k_{FL}C_{DL}'')^{2}}, \qquad (8)$$

where C'_{DL} and C''_{DL} are the real and imaginary parts of the complex capacitance of DL, respectively. The initial capacitance $C'_{in,Total}$ and the slope k_{Total} of the total circuit can be obtained if C'_{Total} is linearly dependent on C''_{Total} .

In the first segment of the parametric study, the parameters of C'_{in} , $_{FL}$ = 200 pF, k_{FL} = 3 $\pi/4$, and C''_{FL} between 5 and 6.5 pF are used to generate the data of FL following Eq. (5). Figures 3(a) and 3(b) exhibit the calculated results on the relationship between C'_{Total} and C''_{Total} with different C'_{DL} and $\tan \delta_{DL}$, where $\tan \delta_{DL}$ is the loss tangent of DL and defined as the ratio of the imaginary and real parts of the complex permittivity. It is clear that C'_{Total} varies linearly with C''_{Total} . The derived $C'_{in,Total}$ and k_{Total} are displayed in Figs. 3(c) and 3(d), respectively. C'_{in} , Total increases with C'_{DL} , but the variation of $C'_{in,Total}$ with $\tan \delta_{DL}$ is not monotonous. As shown in Fig. 3(c), when $\tan \delta_{DL}$ is lower, $C'_{in,Total}$ decreases with tan δ_{DL} increasing. The trend is inverse when tan δ_{DL} is higher. In contrast with $C'_{in,Total}$, different changes of k_{Total} can be observed from Figs. 3(b) and 3(d). k_{Total} decreases with tan δ_{DL} increasing but changes non-monotonically with C'_{DL} . When $\tan \delta_{DL} < 0.03$, k_{Total} is higher than the value of k_{FL} (3 π /4) and decreases gradually toward $3\pi/4$ with C'_{DL} increasing, which corresponds to the DL thickness decreasing. Nevertheless, k_{Total} is lower than $3\pi/4$ when tan $\delta_{DL} \geq 0.03$ and closes up to $3\pi/4$ with C'_{DL} increasing. The results of the parametric study indicate that the value 0.03 for tan δ_{DL} is a demarcation point. According to the parameters that are used for the parametric study, the loss of FL is about 0.03. This means that k_{Total} is not only determined by the capacitance of FL and DL but also related to the relative value of the losses of FL and DL. Therefore, it can be inferred that the reported values of k may be influenced by the interfacial layer between the ferroelectrics and the electrodes.^{7,14,15}

In the second segment of the parametric study, the influence of DL on the frequency dependence of the composite structure is studied. The experimental C-OSC results of BLT film are used as the parameters of FL. DL is represented by a parallel connection unit composed of a resistor and a constant phase element (CPE). The complex capacitance of this parallel unit can be expressed as¹⁰

$$C_{R-CPE}^{*} = Y_0 \omega^{n-1} \sin(n\pi/2) - j \frac{1 + RY_0 \omega^n \cos(n\pi/2)}{R\omega}, \quad (9)$$

where *R* is the resistance, Y_0 represents the pseudo-capacitance, and *n* is the exponent. C^*_{R-CPE} with different *R*, Y_0 , *n*, and ω is calculated by

Eq. (9), and the corresponding C'_{DL} and $\tan \delta_{DL}$ are used to research the effects of DL on the NDR parameters.

The calculated frequency spectra are exhibited in Fig. 4. C'_{DL} and tan δ_{DL} of R-CPE with different R are shown in Figs. 4(a) and 4(b), which represent the case that a ferroelectric capacitor connects in series with the elements that possess same capacitance and different leakage. Neither the frequency nor the resistance affects C'_{DL} tan δ_{DL} decreases with the frequency increasing and correlates with R negatively. The spectra of $C'_{in,Total}$ and k_{Total} are displayed in Figs. 4(c) and 4(d). Compared with BLT, the existence of R-CPE leads to the decrease in $C'_{in,Total}$. In Fig. 4(c), the influence of R mainly appears in the lowfrequency region. It is related to the significant increase in tan δ_{DL} in Fig. 4(b). Higher loss in the low-frequency region results in the larger reduction of $C'_{in, Total}$, especially for the case of $R = 1 \times 10^7 \Omega$. Most of the calculated k_{Total} are higher than the measured k of BLT. It is because the corresponding $\tan \delta_{DL}$ (in the full frequency region of $R = 4 \times 10^7$, 1×10^8 , and $1 \times 10^9 \Omega$ and the high-frequency region of $R = 1 \times 10^7 \Omega$) of the R-CPE is lower than 0.03. On the contrary, for $R = 1 \times 10^7 \Omega$, tan δ_{DL} in the low-frequency region is higher than 0.03. It leads to lower k_{Total} than that of BLT. For the R-CPE, the lower value of R means the leakage current in DL is higher. The calculated results in Fig. 4(d) reveal that the influences of the leakage in DL on k_{Total} cannot be ignored until the resistance of DL is higher than $10^8 \Omega$.

The influences of Y_0 are shown in Figs. 4(e)–4(h). C'_{DL} does not change with frequency but varies with Y_0 because the greater Y_0 causes the higher equivalent capacitance. Moreover, tan δ_{DL} is negatively correlated with Y_0 although it still decreases when the frequency increases.

These effects are opposite to those of *R*. The curves of $C'_{in,Total}$ and k_{Total} are shown in Figs. 4(g) and 4(h), respectively. The influence of the lower tan δ_{DL} on $C'_{in,Total}$ and k_{Total} mainly follows the trend in the region where tan δ_{DL} is less than 0.03 in the mapping diagrams in Fig. 3. In Fig. 4(g), the $C'_{in,Total}$ curves are beneath that of the pure BLT and drop further with C'_{DL} decreasing. However, the k_{Total} curves in Fig. 4(h) are above that of the pure BLT and further increase with the decrease in C'_{DL} . For n = 1, the CPE is equivalent to a capacitor. The calculated results indicate that the influence produced by an interfacial layer with excellent insulation can be regarded as that produced by an ideal capacitor. The linear relationship between the reciprocals of $C'_{in,Total}$ and that of Y_0 in the inset of Fig. 4(g) manifests that the result follows the model that an ideal capacitor connects with the BLT in series.

The effects of *n* are displayed in Figs. 4(i)–4(l). Different from the case of n = 1 in which C'_{DL} does not change with the frequency, C'_{DL} decreases with the increase in frequency when n < 1. The smaller the value of *n*, the greater the curves of C'_{DL} decline. The relationship between C'_{DL} and frequency is linear in double logarithmic coordinates, as shown in Fig. 4(i). tan δ_{DL} decreases with the frequency increasing. The decrease in *n* results in the rapid rise of tan δ_{DL} , as exhibited in Fig. 4(j). The excessively higher tan δ_{DL} means that the ac conductance in CPE is greater because tan δ_{DL} is related to C'_{DL} , which is proportional to the real part of the complex admittance. Therefore, the decline of C'_{DL} and the increase in tan δ_{DL} , as shown in Fig. 4(k). According to the curves in Fig. 4(j), except for the case n = 1, the values of tan δ_{DL} are higher than 0.03 in the cases n = 0.95, 0.9, and 0.7.



FIG. 4. The calculated frequency spectra of (a), (c), and (i) C'_{DL} , (b), (f), and (j) tan δ_{DL} of R-CPE, (c), (g), and (k) $C'_{in, Total}$, and (d), (h), and (l) k_{Total} of the composite system. (a)–(d) $Y_0 = 3 \times 10^{-10}$ S and n = 1 with different R. (e)–(h) $R = 1 \times 10^9 \Omega$ and n = 1 with different Y_0 . (i)–(l) $R = 1 \times 10^9 \Omega$ and $Y_0 = 3 \times 10^{-10}$ S with different n. The inset of (g) shows $1/C'_{in, Total}$ vs $1/Y_0$.



FIG. 5. (a) The capacitance and loss spectra of BLT and Al₂O₃ films. (b) The frequency dependence of $C'_{in,Total}$ and k_{Total} of BLT–AO₁₀₀ and BLT–AO₅₀₀ from the calculated complex capacitance. The frequency spectra of the total equivalent (c) $\varepsilon'_{in,Total}$ and (d) α_{Total} are derived from the calculated C-OSC results. The experimental results of BLT are also plotted in (b)–(d).

Therefore, the values of k_{Total} are lower than that of BLT and decrease significantly with *n*, except for the case n = 1, as shown in Fig. 4(1).

Finally, the experimental data and the model are used conjunctly to research the NDR in our BLT-Al2O3 system. According to the calculated results, k_{Total} of BLT-AO₁₀₀ and BLT-AO₅₀₀ lower than that of BLT in Fig. 2(d) implies that the loss of the Al_2O_3 layer is higher than that of the BLT layer in the composite system. In general, the experimental data for a dead layer are not available. In our work, the experimental data of the Al₂O₃ films on Pt electrode are used. The experimental capacitance and loss spectra of BLT, AO₁₀₀, and AO₅₀₀ are displayed in Fig. 5(a). The loss of BLT is smaller than 0.025, and those of AO₁₀₀ and AO₅₀₀ are higher than 0.05. The capacitance and loss of AO_{100} and AO_{500} in Fig. 5(a) are fitted by Eq. (9), and the fitting results of *R*, *Y*₀, and *n* are $3.93 \times 10^7 \Omega$, 5.20×10^{-10} S, and 0.95, and $2.51\times10^8~\Omega,~1.03\times10^{-10}$ S, and 0.96 for AO_{100} and AO_{500}, respectively. These fitting results and the experimental data of BLT are taken into Eq. (8). The calculated $C'_{in, Total}$ and k_{Total} from the linear relationship between C' and C'' are shown in Fig. 5(b). The results in Fig. 5(b) are similar to the experimental results in Fig. 2(d). Both $C'_{in,Total}$ and k_{Total} decrease with the Al₂O₃ thickness increasing. Based on the calculated results, the smaller k_{Total} and $C'_{in,Total}$ than those of BLT is decided by the capacitance and the loss of the Al2O3, especially for the case of AO_{100} which resistance is lower than $10^8 \Omega$. The effects of DL on α_{Total} are also calculated by our method. The frequency dependence of $\varepsilon'_{in,total}$ and α_{Total} obtained from the calculated C-OSC data is plotted in Figs. 5(c) and 5(d), respectively. The curves exhibit the same trends as the experimental results shown in Fig. 2(d). The fitting results of the calculated curves according to Eq. (4) are also listed in Table I. Compared with BLT-AO₁₀₀, the difference between the experimental and the calculated results for BLT-AO₅₀₀ is higher, which can be ascribed to the smaller capacitance of the AO₅₀₀. The consistency between the experimental and calculated results indicates that our method is appropriate for the studies on NDR in the FE-DE composite system. According to the experimental and calculated results, the NDR parameters in the FE-DE composite system are affected by both the capacitance and the loss of DL and FL, especially when DL is very thin. This work will be significant for the devices based on ultra-thin ferroelectric films and composite systems.

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AUTHOR DECLARATIONS

Conflict of Interest

The authors have no conflicts to disclose.

Author Contributions

Zixin Cao: Data curation (equal); Investigation (equal); Visualization (equal); Writing – original draft (equal). Yawei Li: Conceptualization (equal); Data curation (equal); Investigation (equal); Methodology (equal); Writing – original draft (equal). Liyan Shang: Data curation (equal); Investigation (equal); Writing – original draft (equal). Kai Jiang: Investigation (equal); Writing – original draft (equal). Kai Jiang: Investigation (equal); Methodology (equal); Writing – review & editing (equal). Jinzhong Zhang: Conceptualization (equal); Visualization (equal); Writing – review & editing (equal). Liangqing Zhu: Data curation (equal); Writing – review & editing (equal). Zhigao Hu: Funding acquisition (equal); Methodology (equal); Supervision (equal); Writing – review & editing (equal). Conceptualization (equal); Writing – review & editing (equal). Junhao Chu: Conceptualization (equal); Supervision (equal); Writing – review & editing (equal).

DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding authors upon reasonable request.

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Applied Physics Letters

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